



## Investigation of Ti—B nanoheterofullerenes evolved from C<sub>20</sub> nanocage through DFT

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### ABSTRACT

In present computational survey, we are compared and contrasted electronic effects of C<sub>20</sub> and its C<sub>20-2n</sub>Ti<sub>n</sub>B<sub>n</sub> heterofullerenic derivatives with n = 1-5, at density functional theory (DFT). All nanocages are true minima and none of them collapses to deformed open cage as segregated nanostructure. Isolating five single hetero bonds of Ti—B *via* either one double bond of C=C or one carbon atom is a suitable method for reaching highly substituted stable heterofullerene *i.e.* pen-shell C<sub>10</sub>Ti<sub>5</sub>B<sub>5</sub> since it prevents from weak homo bonding of Ti—Ti and B—B. The C<sub>7</sub>-C<sub>18</sub>Ti<sub>1</sub>B<sub>1</sub> heterofullerene can avoid from the most strained directly fused five-pentagon configuration, but its open-shell electronic structure with highly pyramidalized titanium atom (126.59° *i.e.* 3-4 times relation to C<sub>20</sub>) may render it too reactive to be observed under typical experimental conditions. The calculated binding energy (*B.E.*) shows C<sub>10</sub>Ti<sub>5</sub>B<sub>5</sub> as the most stable heterofullerene. Contrary on *B.E.*, the absolute heat of atomization of heterofullerenes decreases as number of substituting Ti—B unit (*n*) increases. Hence, substitution effect on binding energy is more significant than heat of atomization. Compared to the previously reported material such Ti-decorated B<sub>38</sub> nanofullerene as a suitable hydrogen storage with *B.E.* of 5.67 eV/atom, our studied C<sub>20-2n</sub>Ti<sub>n</sub>B<sub>n</sub> heterofullerenes show the higher *B.E.* with range of 12.25 to 38.03 eV/atom, the higher stability and the higher capacity for hydrogen storage. Interestingly, C<sub>18</sub>Ti<sub>1</sub>B<sub>1</sub> heterofullerene must be not only kinetically stable against electronic excitations but also based on natural bond orbital (NBO) analysis, the highest charge transfer (CT) is take placed from π<sub>C=C</sub> bonds to the neighboring LP\*<sub>B</sub>, LP\*<sub>Ti</sub> and σ\*<sub>Ti-B</sub> anti-bonds of it. Also, the exceptionally large value of nucleus-independent chemical shifts (NICS) is found for C<sub>18</sub>Ti<sub>1</sub>B<sub>1</sub> compared to the archetypically aromatic benzene and [2.2]paracyclophane molecules.

### 1. Introduction

Substituting for carbon atoms in nanostructures with different heteroatoms has been the object of a wealth of experimental and computational studies aimed at unraveling their geometrical and electronic properties for example polymorphism, extraordinary hardness, band gap, chemical reactivity, *etc.* [1]. The special properties of nanostructured synthetic carbon allotropes can be related to the curved π-conjugation within their

curved hexagonal, pentagonal and hybrid hexagon-pentagon carbon-carbon networks [2]. The curvature leads to a loss of spatial overlap of the atomic p orbitals and a shift in hybridization of the carbon atoms from the sp<sup>2</sup> of graphite to a state intermediate between sp<sup>2</sup> and sp<sup>3</sup> [3].

Boron-substituted fullerene is appropriate as a substrate since it consists of light elements and has the capability to bind with metal atoms such as Ti, Zr and U [4,5]. After B<sub>80</sub> fullerene [6], different types of B-doped

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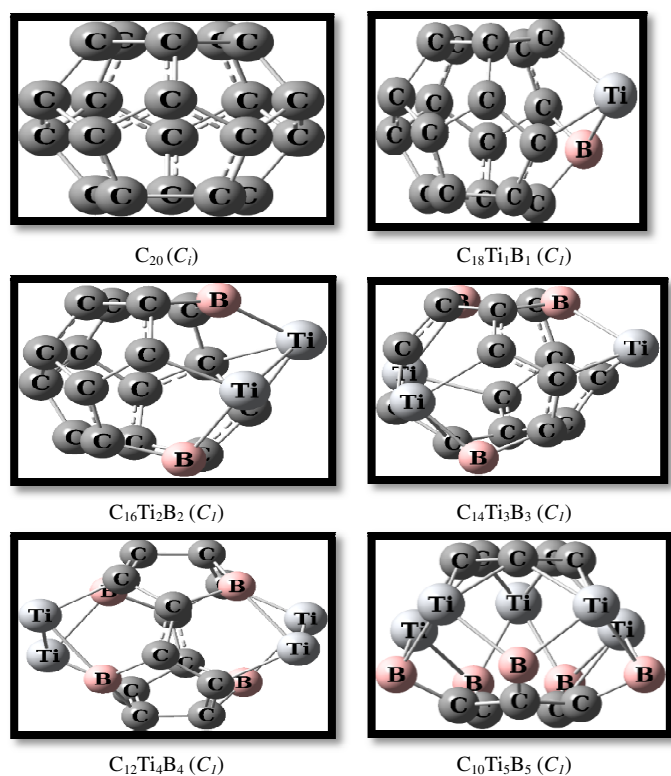
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fullerenes (from  $B_{32}$  to  $B_{80}$ ) have been suggested [4-8]. The alkali and alkaline-earth decorated boron fullerenes ( $B_{38}$  and  $B_{80}$ ) have been commonly suggested as candidates for hydrogen storage [4-8]. The Ti-decorated fullerenes ( $B_{38}$ ,  $B_{40}$  and  $C_{24}B_{24}$ ) along with other nanostructures (Ti-coated graphene nanoribbons) also have been widely suggested as potential hydrogen storage materials [4-8]. All of the research mentioned above encourages us to examine whether the transition metal Ti atom decorated  $B_{38}$  fullerene is an efficient hydrogen storage media. There is little research on hydrogen storage and other properties of nanostructures of the form  $C_{20-2n}Ti_nB_n$  [9-13]. Since the bulk boron cannot be found in nature, the design and synthesis of bulk boron allotropes and boron-substituted fullerenes still keeps challenging to theoretical and experimental chemists. Boron fullerenes are seen as efficient hydrogen storage media since they are light-weight and have the capability to bind with metalead atoms. Since isolated transition metal has the ability to bind a certain number of hydrogens, theoretical simulations on hydrogen adsorption by Ti—B heterofullerenes are reported as organotitanium compounds that we hope provide useful information. Thus, the aim of this paper is to survey hydrogen storage capacity of  $C_{20-2n}Ti_nB_n$  heterofullerenes ( $n = 1-5$ ) by DFT (Figure 1)

The Ti atoms can be substituted to carbon atoms of the  $C_{20-n}B_n$  heterofullerenes ( $n = 1-5$ ) due to the CT from the Ti atom(s) to boron and carbon atoms. Our results indicate that a strong binding among Ti atoms and the substituted boron of heterofullerenes can be efficiently prevents from segregation of  $C_{20-2n}Ti_nB_n$  nanocages.

## 2. Computational Methods

In this survey, full geometry optimization and harmonic frequency calculations of  $C_{20}$  and  $C_{20-2n}Ti_nB_n$  heterofullerenes are carried out without any symmetry constraints using hybrid functional B3LYP/6-311+G\* [14] and the GAMESS program [15]. To reach more accurate energetic data, single point calculations for closed-shell multiplicity ( $n = 0, 2$  and  $4$ ) are accomplished at restricted spin symmetry methods of B3LYP/6-311++G\*\*, M06-2X/6-311++G\*\* [16] and B3LYP/AUG-cc-pVTZ based on the B3LYP/6-311+G\* and M06-2X/6-311+G\* geometries, respectively [17]. Structures with open-shell multiplicity ( $n = 1, 3$  and  $5$ ) are optimized using unrestricted spin symmetry methods including UB3LYP/6-311++G\*\*, UM06-2X/6-311++G\*\* and UB3LYP/AUG-cc-pVTZ based on UB3LYP/6-311+G\* and UM06-2X/6-311+G\* geometries, respectively. As a thermodynamic stability criterion of nanocages, binding energies are calculated according to the following expression:  $B.E. = [(20-2n)E_C + nE_{Ti} + nE_B - E_{tot}] / 20$  where  $E_C$ ,  $E_{Ti}$ ,  $E_B$ ,  $E_{tot}$  is the energy of carbon, titanium, boron atoms and total energy of the fullerene, respectively. The  $E_C$ ,  $E_{Ti}$  and  $E_B$  values are guessed to be  $-37.86$ ,  $-844.42$  and  $-24.66$  hartree, respectively, at the B3LYP/6-311+G\*. As a kinetic stability criterion of nanocages, the NBO analysis on optimized structures is achieved, at B3LYP/AUG-cc-pVTZ, M06-2X/6-311++G\*\* and B3PW91/6-311++G\*\* [18, 19]. The NICS parameter is calculated at the centers of nanocages [20]. In addition, heat of atomization ( $\Delta H_{at}$ ) is accessible *via* the following expression:  $\Delta H_{at}(C_{20-2n}Ti_nB_n) = \Delta H_f(C_{20-2n}Ti_nB_n) - (20-2n)\Delta H_f(C) - n\Delta H_f(Ti) - n\Delta H_f(B)$  where  $\Delta H_f$  symbolizes the heat of formation. Then  $\Delta H_f(C_{20-2n}Ti_nB_n) = E_{tot} - (20-2n)E_C - nE_{Ti} - nE_B$ , where  $E_{tot}$ ,  $E_C$ ,  $E_{Ti}$  and  $E_B$ , is total energy of heterofullerene, energy of carbon, titanium and boron atoms, respectively. The experimental values for  $H_f(C)$ ,  $H_f(Ti)$  and  $H_f(B)$ , are  $171.37$ ,  $111.85$  and  $77.92$  kcal/mol, respectively [21]. The nucleophilicity index,  $N$ , is calculated as  $N = E_{HOMO(Nu)} - E_{HOMO(TCNE)}$ ; (tetracyanoethylene (TCNE) is preferred as the reference) [22]. The global electrophilicity,  $\omega$  is calculated *via* the following expression,  $\omega = (\mu^2 / 2\eta)$ , where  $\mu$  is the chemical

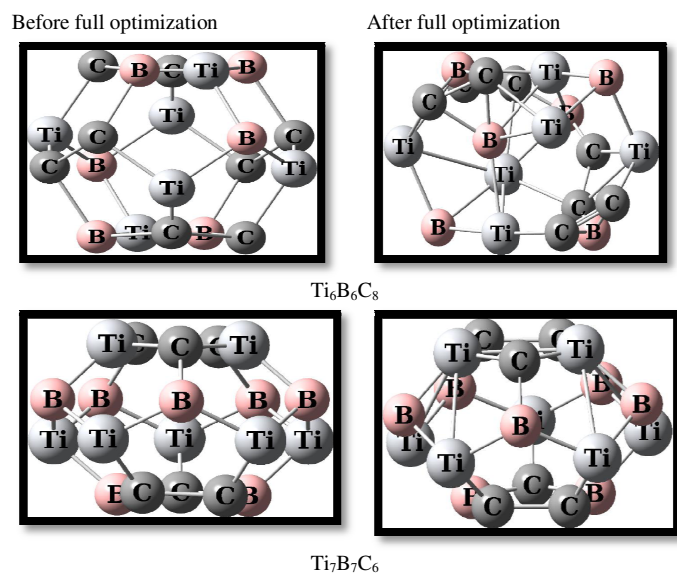


**Fig. 1.** Full optimized structures of  $C_{20}$  and  $C_{20-2n}Ti_nB_n$  heterofullerenes ( $n = 1-5$ ) along with their corresponding symmetries, at B3LYP/AUG-cc-pVTZ.

potential ( $\mu = (E_{\text{HOMO}} + E_{\text{LUMO}}) / 2$ ) and  $\eta$  is the global hardness ( $\eta = (E_{\text{LUMO}} - E_{\text{HOMO}}) / 2$ ) [22].

### 3. Results and discussion

Before the optimization of the scrutinized structures such as  $\text{Ti}_6\text{B}_6\text{C}_8$  and  $\text{Ti}_7\text{B}_7\text{C}_6$  we completed conformer distribution analysis by Spartan 10 program and attained their relative energies along with their stabilities. Then, full optimization done without any symmetry constraints using the same methods and basis sets under study in this work (Figure 2).



**Fig. 2.** Full optimized structures of  $\text{Ti}_6\text{B}_6\text{C}_8$  and  $\text{Ti}_7\text{B}_7\text{C}_6$  heterofullerenes, at B3LYP/AUG-cc-pVTZ.

Apparently, the fullerene structures of  $\text{Ti}_6\text{B}_6\text{C}_8$  and  $\text{Ti}_7\text{B}_7\text{C}_6$  lose their uniformity due to the higher covalent radii of titanium (1.40 Å) compared to carbon (0.70 Å) and boron (0.85 Å) also shorting of Ti—Ti distances in equatorial and cap positions, *i.e.*, increasing the number of substituted Ti more than five, moves the titanium heteroatoms from their natural positions to formation intra-molecular bonding. Hence, these heterofullerenes deform their cage at the Ti—Ti bonds which make them appear cubic-like with the longer Ti—Ti bond lengths (2.872 Å). This phenomenon is occurred in order to reduce their curvature effects and to release their strain energy. The formation of these clusters is quite complicated due to the possibility of high-spin states. According to effort which has been made for global optimization, argument seems rationale for adopting the specific cage structures shown in Figure 1.

#### 3.1. Structural parameters, energy and stability

The average C—C bond distance of  $\text{C}_{20}$  is from 1.444 to 1.537 Å which enclosures among the bond

lengths of ethylene and ethane molecules (1.35, 1.54 Å, respectively) (Tables 1 and 2).

**Table 1.** The range of bond lengths (Å) calculated for  $\text{C}_{20}$  and its  $\text{C}_{20-2n}\text{Ti}_n\text{B}_n$  heterofullerenes, at B3LYP/AUG-cc-pVTZ.

Species	C=C	Ti—C	B—C	Ti—B	Ti—Ti
$\text{C}_{20}$	1.444-1.537	-	-	-	-
$\text{C}_{18}\text{Ti}_1\text{B}_1$	1.423-1.479	2.086	1.571	2.187	-
$\text{C}_{16}\text{Ti}_2\text{B}_2$	1.426-1.471	1.924-2.207	1.527-1.646	2.042-2.151	2.646
$\text{C}_{14}\text{Ti}_3\text{B}_3$	1.371-1.472	1.902-1.912	1.524-1.691	2.002-2.007	-
$\text{C}_{12}\text{Ti}_4\text{B}_4$	1.399-1.469	2.064	1.556-1.658	2.247	2.284
$\text{C}_{10}\text{Ti}_5\text{B}_5$	1.450-1.498	2.086-2.133	1.495-1.527	2.093-2.203	-

**Table 2.** The range of bond angles of C—C—C, C—Ti—C, C—B—C (degree) and  $\theta_{\text{Ti}}$ ,  $\theta_{\text{B}}$  pyramidalization angles (degree) calculated for  $\text{C}_{20}$  and its  $\text{C}_{20-2n}\text{Ti}_n\text{B}_n$  heterofullerenes, at B3LYP/AUG-cc-pVTZ. The pyramidalization angle of  $\text{C}_{20}$  is  $\theta_{\text{C}}$ .

Species	C—C—C	C—Ti—C	C—B—C	$\theta_{\text{Ti}}$	$\theta_{\text{B}}$
$\text{C}_{20}$	103.99-111.31	-	-	-	-
$\text{C}_{18}\text{Ti}_1\text{B}_1$	103.35-114.23	79.64	100.72	126.59	36.22
$\text{C}_{16}\text{Ti}_2\text{B}_2$	105.23-117.92	-	101.61-102.64	62.15	52.23
$\text{C}_{14}\text{Ti}_3\text{B}_3$	107.70-112.89	83.81-87.40	80.23	61.24	51.16
$\text{C}_{12}\text{Ti}_4\text{B}_4$	109.23-112.34	-	97.96	107.35-131.25	38.06-46.35
$\text{C}_{10}\text{Ti}_5\text{B}_5$	106.21-109.67	39.86-40.89	-	-	-

The substituted doping increases bond lengths of Ti—C, B—C, Ti—B and Ti—Ti to 1.495 Å – 2.646 Å; which are closed to the sum of covalent radii of carbon (0.70 Å), titanium (1.40 Å) and boron (0.85 Å); and decreases bond lengths of C=C to 0.039 Å – 0.073 Å. Also, the bond angles of C—Ti—C and C—B—C are diminished in the range of 39.86° – 87.40° and 80.23° – 102.64°, respectively, vs. the range of C—C—C bond angles of C<sub>20</sub> from 103.99° to 111.31° (Table 2). The pyramidalization angle on C, Ti and B atoms of C<sub>20</sub> and C<sub>18</sub>Ti<sub>1</sub>B<sub>1</sub> is measured as  $\theta_C = 360 - [(C-C-C)_1 + (C-C-C)_2 + (C-C-C)_3]$ ,  $\theta_{Ti} = 360 - [(C-Ti-C)_1 + (C-Ti-B)_2 + (B-Ti-C)_3]$  and  $\theta_B = 360 - [(C-B-C)_1 + (C-B-Ti)_2 + (Ti-B-C)_3]$ , respectively [23]. Since titanium and boron heteroatoms adopt *sp*<sup>3</sup> hybridization more than *sp*<sup>2</sup>, these substitutions display higher pyramidalization angle of  $\theta_{Ti}$  and  $\theta_B$  than  $\theta_C$ , and deform their cage around Ti—B bonds. Also, the absolute value of  $E_{tot}$  is increased as n of heterofullerenes is increased and the smallest vibrational frequency ( $\nu_{min}$ ) emerges true minimum for them (Table 3).

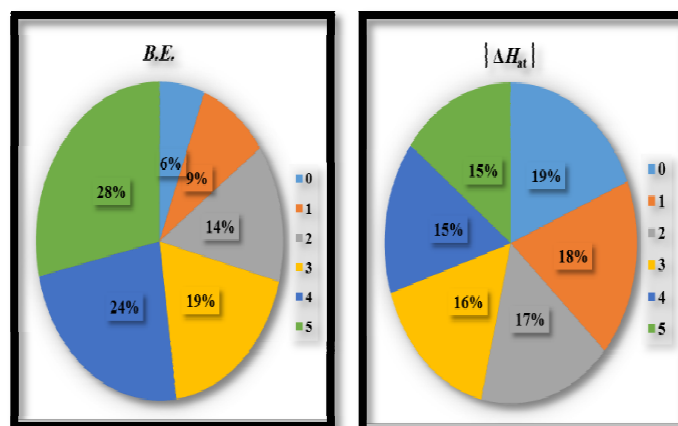
**Table 3.** The first thermodynamic parameters including total energies ( $E_{tot}$  in *a.u.*), smallest vibrational frequencies ( $\nu_{min}$  in cm<sup>-1</sup>), zero-point vibrational energies (ZPVE in kcal/mol) and binding energies (*B.E.* in eV/atom) calculated for C<sub>20</sub> and its C<sub>20-2n</sub>Ti<sub>n</sub>B<sub>n</sub> derivatives with n = 1-5.

Species	$E_{tot}^{a, (b), c, [d]}$	$\nu_{min}^e$	ZPVE <sup>e</sup>	<i>B.E.</i> <sup>d</sup>
C <sub>20</sub>	-761.60234 (-761.18425) -761.18350 [-761.67038]	82.23	70.84	7.83
C <sub>18</sub> Ti <sub>1</sub> B <sub>1</sub>	-1559.563072 (-) 1559.207676) -1559.44292 [- 1559.852027]	167.41	66.63	12.25
C <sub>16</sub> Ti <sub>2</sub> B <sub>2</sub>	-2357.639009 (-) 2357.316075) -2357.526294 [- 2357.997112]	144.46	61.64	18.67
C <sub>14</sub> Ti <sub>3</sub> B <sub>3</sub>	-3155.677780 (-) 3155.363587) -3155.555512 [- 3156.095962]	106.52	53.84	25.03

C <sub>12</sub> Ti <sub>4</sub> B <sub>4</sub>	-3953.747407 (-) 3953.453624) -3953.626146 [- 3954.229619]	62.39	51.72	31.44
C <sub>10</sub> Ti <sub>5</sub> B <sub>5</sub>	-4751.95320 (-) 4751.704444) -4751.83972 [- 4751.900421]	24.24	45.66	38.03

At <sup>a</sup>B3LYP/6-311++G\*\*, <sup>b</sup>M06-2X/6-311++G\*\*, <sup>c</sup>B3PW91/6-311++G\*\*, <sup>d</sup>B3LYP/AUG-cc-pVTZ, and <sup>e</sup>B3LYP/6-311+G\*.

Probing the mixed-cage geometries show the reduction of  $\theta_B$  pyramidalization angles to compensate for the bigger  $\theta_{Ti}$  pyramidalization angles *and vice versa* (Figure 3).

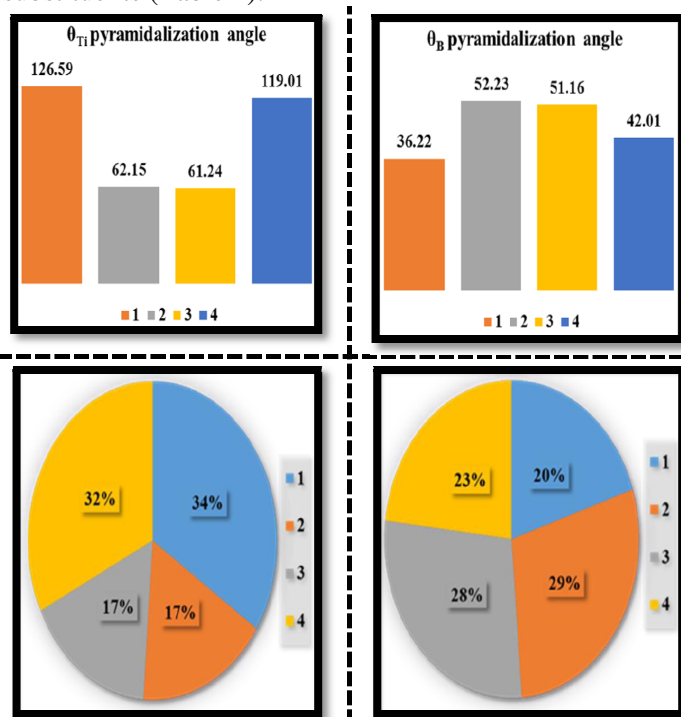


**Fig. 3.** Comparison of calculated binding energy (*B.E.* in eV/atom), and absolute heat of atomization ( $|\Delta H_{at}|$  in kcal/mol) for C<sub>20</sub> cage (n = 0) and C<sub>20-2n</sub>Ti<sub>n</sub>B<sub>n</sub> heterofullerenes (n = 1-5), at B3LYP/6-311+G\*.

Substituent effects of titanium and boron on degrees of deformation show differently percentage, as the  $\theta_B$  pyramidalization angles are smaller than the  $\theta_{Ti}$  pyramidalization angles. It seems that, from C<sub>18</sub>Ti<sub>1</sub>B<sub>1</sub> to C<sub>12</sub>Ti<sub>4</sub>B<sub>4</sub>; the  $\theta_B$  angle, tendency of curvature and the released strain energy in the boron sites decreases to compensate for increasing of  $\theta_{Ti}$  angle, tendency of curvature and the released strain energy in the titanium sites *and vice versa*. Furthermore, due to relief of strain, the bond angles and the  $\theta_{Ti}$ ,  $\theta_B$  pyramidalization angles are not comparable to each other. Due to the natural *sp*<sup>2</sup> hybridization of boron, and in order to accommodate the smaller B atoms in the cage structure, all  $\theta_B$



pyramidalization angles are smaller than  $100^\circ$  while  $\theta_{\text{Ti}}$  pyramidalization angles are either smaller or bigger than  $100^\circ$ . The  $\text{C}_{18}\text{Ti}_1\text{B}_1$ ,  $\text{C}_{16}\text{Ti}_2\text{B}_2$  and  $\text{C}_{14}\text{Ti}_3\text{B}_3$  species show the  $\nu_{\text{min}}$  values higher than  $\text{C}_{20}$  (167.41, 144.46 and  $106.52\text{ cm}^{-1}$ , respectively, vs.  $82.23\text{ cm}^{-1}$ ) [24]. Interestingly, there is a rather good consistency between the absolute value of  $E_{\text{tot}}$ , zero-point vibrational energies (ZPVE) and *B.E.* while *n* increases. Thus, substituting increases *B.E.* of the scrutinized heterofullerenes (the range of 12.25 eV/atom for  $\text{C}_{18}\text{Ti}_1\text{B}_1$  to 38.03 eV/atom for  $\text{C}_{10}\text{Ti}_5\text{B}_5$ ) compared to the unsubstituted cage (7.83 eV/atom) (Figure 4). The frontier molecular orbital's (FMO's) energy splitting has been applied as a criterion of kinetic stability of fullerenes and other nanostructures. Here, the band gap of structures is differently depending on *n* and topology of Ti and B substituents (Table 4).



**Fig. 4.** Substituent effects of titanium (left) and boron (right) on stability of  $\text{C}_{20-2n}\text{Ti}_n\text{B}_n$  heterofullerenes *via* comparison of the resulted pyramidalization angles (degree) and *n*, at B3LYP/AUGcc-pVTZ.

**Table 4.** The first kinetic parameters including the frontier molecular orbital energies ( $E_{\text{HOMO}}$  and  $E_{\text{LUMO}}$  in *a.u.*), along with their band gaps ( $\Delta E_{\text{HOMO-LUMO}}$  in eV) calculated for  $\text{C}_{20}$  and its  $\text{C}_{20-2n}\text{Ti}_n\text{B}_n$  heterofullerenic derivatives with *n* = 1-5.

Species	$E_{\text{HOMO}}^{\text{a, (b), [c]}}$	$E_{\text{LUMO}}^{\text{a, (b), [c]}}$	$\Delta E_{\text{HOMO-LUMO}}^{\text{a, (b), [c]}}$
$\text{C}_{20}$	-0.20134, (-0.21055) [-0.20455]	-0.13054 (-0.12567) [-0.13351]	1.93 (2.31) [1.93]

$\text{C}_{18}\text{Ti}_1\text{B}_1$	-0.19888 (-0.19898) [-0.20178]	-0.11656 (-0.10724) [-0.11726]	2.24 (2.50) [2.30]
$\text{C}_{16}\text{Ti}_2\text{B}_2$	-0.18848 (-0.19307) [-0.19002]	-0.12314 (-0.11716) [-0.12221]	1.78 (2.07) [1.85]
$\text{C}_{14}\text{Ti}_3\text{B}_3$	-0.17193 (-0.16181) [-0.16149]	-0.09496 (-0.08818) [-0.09599]	2.09 (2.00) [1.78]
$\text{C}_{12}\text{Ti}_4\text{B}_4$	-0.15980 (-0.16426) [-0.16069]	-0.08961 (-0.08456) [-0.08951]	1.91 (2.17) [1.94]
$\text{C}_{10}\text{Ti}_5\text{B}_5$	-0.15961 (-0.16760) [-0.16809]	-0.10512 (-0.10425) [-0.11018]	1.48 (1.72) [1.58]

At <sup>a</sup>B3LYP/AUG-cc-pVTZ, <sup>b</sup>M06-2X/6-311++G\*\*, and <sup>c</sup>B3PW91/6-311+G\* levels of theory.

Based on B3LYP/AUG-cc-pVTZ, M06-2X/6-311++G\*\* and B3PW91/6-311+G\* calculations, the  $\Delta E_{\text{HOMO-LUMO}}$  of  $\text{C}_{20}$  fullerene is estimated 1.93, 2.31 and 1.93 eV, respectively, while the  $\Delta E_{\text{HOMO-LUMO}}$  of  $\text{C}_{18}\text{Ti}_1\text{B}_1$  by substituting of one Ti—B unit is changed to 2.24, 2.50 and 2.30 eV, respectively). The substituted titanium and boron heteroatoms simultaneously stabilize the highly strained  $\text{C}_{20}$  cage by CT. Since the  $\text{C}_{20}$  cage is isolobal but is not isoelectron to the substituted heterofullerenes, the number of electrons and multiplicity of the ground state is varied from a fullerene to another fullerene. Henceforth, the  $\text{C}_{20}$ ,  $\text{C}_{16}\text{Ti}_2\text{B}_2$  and  $\text{C}_{12}\text{Ti}_4\text{B}_4$  structures are electronically closed-shell, while  $\text{C}_{18}\text{Ti}_1\text{B}_1$ ,  $\text{C}_{14}\text{Ti}_3\text{B}_3$  and  $\text{C}_{10}\text{Ti}_5\text{B}_5$  are open-shell.

Moreover, the absolute  $\Delta H_{\text{at}}$  value of heterofullerenes is decreased with increasing *n*. From an energy perspective,  $\text{C}_{10}\text{Ti}_5\text{B}_5$  has the lowest  $|\Delta H_{\text{at}}|$  and it is the most stable species (Table 5).

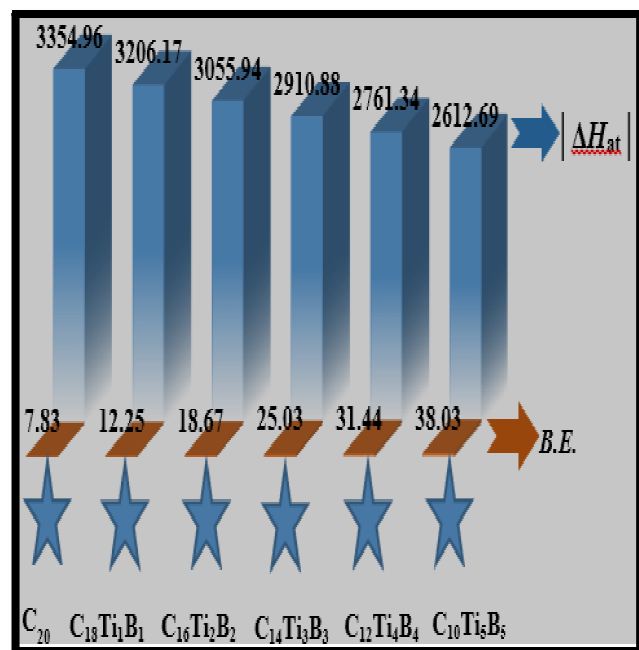
**Table 5.** The second thermodynamic parameters including heat of atomization ( $\Delta H_{\text{at}}$  in kcal/mol), nucleus-independent chemical shifts at cage centers (NICS (0) in ppm), polarity ( $\mu$  in Debye), as well as isotropic polarizability ( $\langle\alpha\rangle$  in *a.u.*) calculated for  $\text{C}_{20}$  and its  $\text{C}_{20-2n}\text{Ti}_n\text{B}_n$  heterofullerenic derivatives with *n* = 1-5.

Species	$\Delta H_{\text{at}}^{\text{a}}$	NICS (0) <sup>b</sup>	NICS (0) <sub>ZZ</sub> <sup>b</sup>	$\mu^{\text{b}}$	$\langle\alpha\rangle^{\text{b}}$
$\text{C}_{20}$	3354.96	-23.81	+3.69	0.00	187.27

$C_{18}Ti_1B_1$	3206.17	-50.83	-71.50	6.02	201.64
$C_{16}Ti_2B_2$	3055.94	-30.33	-59.84	7.52	258.64
$C_{14}Ti_3B_3$	2910.88	-17.02	-30.30	1.86	303.31
$C_{12}Ti_4B_4$	2761.34	+0.93	-10.82	0.00	390.42
$C_{10}Ti_5B_5$	2612.69	+1.93	+0.02	5.60	353.56

At <sup>a</sup>B3LYP/6-311+G\*, <sup>b</sup>B3LYP/AUG-cc-pVTZ.

The  $C_{20}$ ,  $C_{18}Ti_1B_1$ ,  $C_{16}Ti_2B_2$ ,  $C_{14}Ti_3B_3$ ,  $C_{12}Ti_4B_4$  and  $C_{10}Ti_5B_5$  species have 6, 9, 14, 19, 24 and 28% *B.E.* character in stability, respectively, while they have approximately similar percentage of  $|\Delta H_{at}|$  character in stability in the range of 15% to 19% (Figure 4). From the mixing of the two graphs, the share of *B.E.* is different from the share of  $|\Delta H_{at}|$  in thermodynamic stability of  $C_{20}$  and  $C_{20-2n}Ti_nB_n$  nanocages (Figure 5).

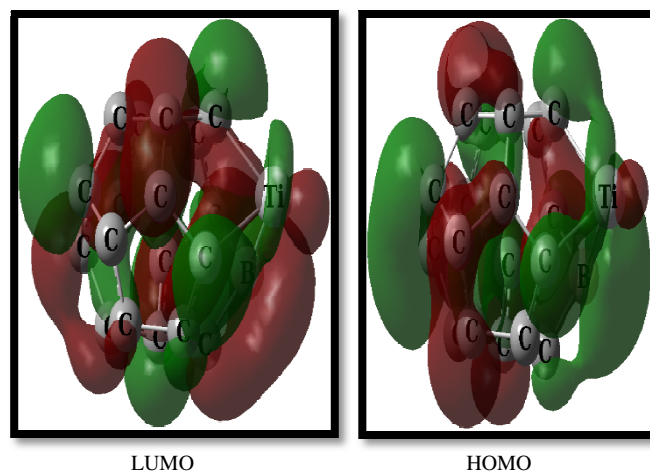


**Fig. 5.** Comparison of stability *via* calculated *B.E.* (eV/atom), and  $|\Delta H_{at}|$  (in kcal/mol) for  $C_{20}$  and its  $C_{20-2n}Ti_nB_n$  derivatives with  $n = 1-5$ , at B3LYP/6-311+G\*.

### 3.2. Aromaticity and charge

Here, we are obtained NICS (0) of  $C_{20}$  -23.8 ppm, which implies its strong aromaticity compared to normal benzene and [2.2]paracyclophane molecules; as

reference models (-8.5 and -9.5 ppm, respectively). Stacking of two rings to give [2.2]paracyclophane decreases NICS in the inner region. However, in the outer region NICS unexpectedly increases. This increase of NICS in the outer region of [2.2]paracyclophane is likely due to the not completely parallel disposition of the two benzene units in this compound. All the  $C_{20-2n}Ti_nB_n$  derivatives show various NICS (0) with range of -50.80 to +1.93 ppm. Thus,  $C_{18}Ti_1B_1$ ,  $C_{16}Ti_2B_2$  and  $C_{14}Ti_3B_3$  nanocages are aromatic, while  $C_{12}Ti_4B_4$  and  $C_{10}Ti_5B_5$  nanocages are anti-aromatic. Also,  $C_{18}Ti_1B_1$  species exhibits the most negative NICS (0) of -50.80 ppm, among six nanocages. The occupancy of two carbon atoms by one titanium and boron atom, improves the overlap between the nearest carbon atoms with atomic orbitals of substituted dopants indicating the enhanced  $\pi$ -aromaticity (Figure 6).



**Fig. 6.** The shapes of HOMO and LUMO orbitals of  $C_{18}Ti_1B_1$ .

Table 5 shows the NICS (0)<sub>zz</sub> values better describe the ring current due to the  $\pi$ -electrons as being the out of plane component of NICS. Moreover, all of the scrutinized nanocages display diverse NICS (0)<sub>zz</sub> with range of -71.50 ppm for  $C_{18}Ti_1B_1$  to +0.02 ppm for  $C_{10}Ti_5B_5$  which imply strong ring current of  $C_{18}Ti_1B_1$ ,  $C_{16}Ti_2B_2$  and  $C_{14}Ti_3B_3$  nanocages (-71.50, -59.84 and -30.30, respectively) compared to  $C_{20}$  (+3.69 ppm). Though the electronegativity difference is considerable (C: 2.5, Ti: 1.54 and B: 2.0),  $\pi$ -delocalization on the ring perimeter is strengthened in  $C_{18}Ti_1B_1$ , while such the ring current is weakened in  $C_{10}Ti_5B_5$  (Figure 7). The Ti heteroatoms of  $C_{20-2n}Ti_nB_n$  bear varying NBO charges from -0.401 (the most negative charge on carbon atom of  $C_{10}Ti_5B_5$ ) to +1.084 (the most positive charge on titanium atom of  $C_{18}Ti_1B_1$  compared to  $C_{20}$  (the range of -0.054 to +0.054), at B3PW91/6-311+G\*\* level (Figure 8).

There is a considerable CT from the substituted electropositive Ti heteroatoms to their adjacent electronegative atoms. Boron substitution induces a slight range of positive charge of +0.110 in  $C_{10}Ti_5B_5$  to +0.402 in  $C_{14}Ti_3B_3$ . Clearly, B and C elements inserted in second row of periodic table, which this fact leads to the highest overlap of empty p orbitals of B heteroatom(s) with  $\pi$  orbitals of C=C bond.

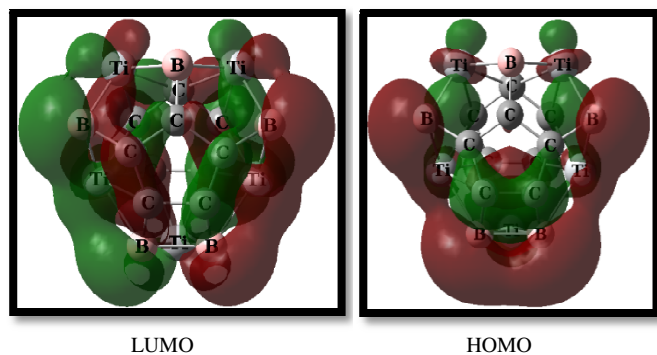
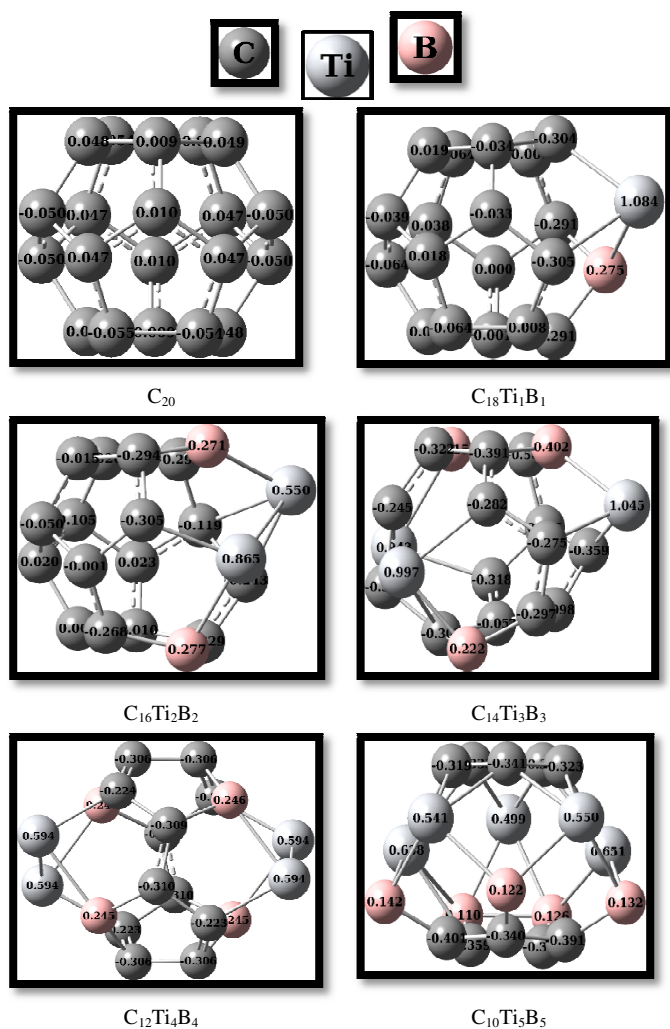
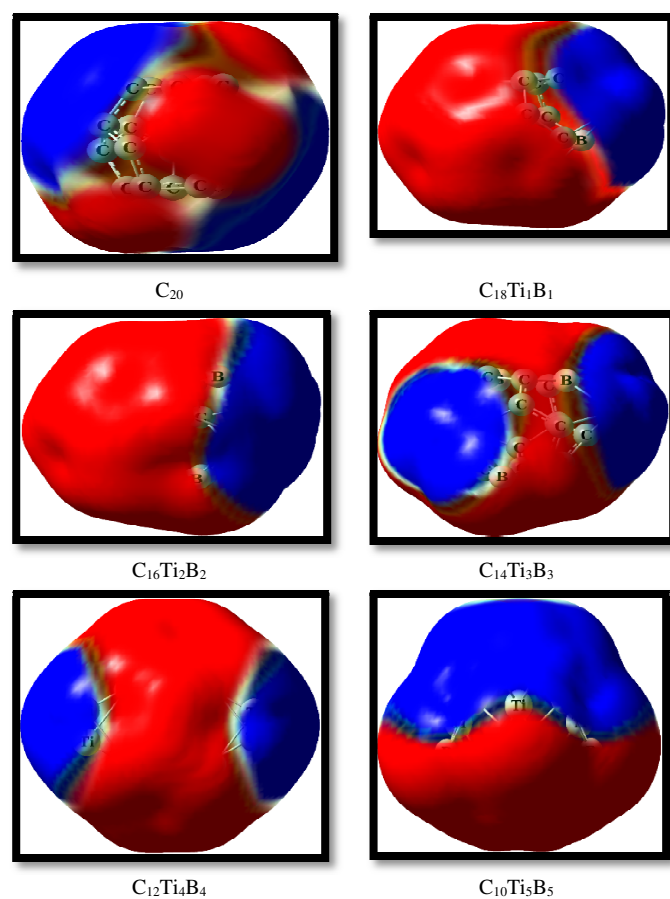


Fig. 7. The shapes of HOMO and LUMO orbitals of  $C_{10}Ti_5B_5$ .



**Fig. 8.** The NBO atomic charge distribution on C, Ti and B atoms calculated for  $C_{20}$  and its  $C_{20-2n}Ti_nB_n$  heterofullerenic derivatives with  $n = 1-5$ , at B3PW91/6-311++G\*\* level of theory.

Also, Ti atom(s) including electron configuration of  $[Ar](4s^2)(3d^2)$ , shares two electron with incomplete p orbital of B atom(s). Hence, there is an overall CT over the surfaces of  $C_{20-2n}Ti_nB_n$  especially  $C_{18}Ti_1B_1$  compared to  $C_{20}$  as the best and as the worst applicant for hydrogen storage, respectively. The MEP plots of the scrutinized heterofullerenes is revealed blue color for titanium sites *via* their positive charge while their negative charge is displayed *via* the red electron cloud in the middle and cap positions for carbon and boron sites. These MEP plots indicate a CT is happened from the Ti heteroatom(s) to the adjacent C and B atoms (Figure 9).



**Fig. 9.** The resulted MEP maps of  $C_{20}$  and its  $C_{20-2n}Ti_nB_n$  heterofullerenic derivatives with  $n = 1-5$ , at B3PW91/6-311++G\*\* level of theory.

Subsequently, Ti—B substituting leads to considerably dipole moment (7.52, 6.02, 5.60 and 1.86 Debye, for  $C_{16}Ti_2B_2$ ,  $C_{18}Ti_1B_1$ ,  $C_{10}Ti_5B_5$  and  $C_{14}Ti_3B_3$  structures, respectively) while  $\mu$  of  $C_{20}$  and  $C_{12}Ti_4B_4$  cages are zero. The average polarizability;  $\langle\alpha\rangle$ ; enlarges from 187.27

*a.u.* for C<sub>20</sub> to 390.42 *a.u.* for C<sub>12</sub>Ti<sub>4</sub>B<sub>4</sub> molecule. Obviously, four Ti—B substituting units to C<sub>20</sub> causes to increasing  $\langle\alpha\rangle$  about two times and this species has high activity to polar-polar interaction.

### 3.3. Reactivity

In this section, the obtained global reactivity of the C<sub>20</sub> and its substituted C<sub>20-2n</sub>Ti<sub>n</sub>B<sub>n</sub> nanocages from FMO's energy is discussed in terms of  $N$ ,  $\omega$ ,  $\mu$ ,  $\eta$ ,  $\chi$ ,  $S$ , along with  $\Delta N_{\max}$ , based on M06-2X/6-311++G\*\* calculations (Table 6).

**Table 6.** The second kinetic parameters including nucleophilicity index ( $N$ ), global electrophilicity ( $\omega$ ), chemical potential ( $\mu$ ), global hardness ( $\eta$ ), electronegativity ( $\chi$ ), global softness ( $S$ ), and also maximum electronic charge ( $\Delta N_{\max}$ ) (all in eV) calculated for C<sub>20</sub> and its C<sub>20-2n</sub>Ti<sub>n</sub>B<sub>n</sub> heterofullerenic derivatives with  $n = 1-5$ , at M06-2X/6-311++G\*\* level of theory.

Species	$N$	$\omega$	$\mu$	$\eta$	$\chi$	$S$	$\Delta N_{\max}$ x
C <sub>20</sub>	3.73	4.53	-4.57	2.31	4.57	0.22	1.98
C <sub>18</sub> Ti <sub>1</sub> B	4.05	3.48	-4.17	2.50	4.17	0.20	1.67
C <sub>16</sub> Ti <sub>2</sub> B <sup>1</sup>	4.21	4.31	-4.22	2.07	4.22	0.24	2.04
C <sub>14</sub> Ti <sub>3</sub> B <sup>2</sup>	5.06	2.89	-3.40	2.00	3.40	0.25	1.70
C <sub>12</sub> Ti <sub>4</sub> B <sup>3</sup>	4.99	2.64	-3.39	2.17	3.39	0.23	1.56
C <sub>10</sub> Ti <sub>5</sub> B <sup>4</sup>	4.90	3.97	-3.70	1.72	3.70	0.29	2.15

Substituting of one to five Ti and B—heteroatoms lead to distinctive variations of  $N$ ,  $\omega$ ,  $\mu$ ,  $\eta$ ,  $\chi$ ,  $S$  and  $\Delta N_{\max}$ . So that, one substituted Ti—B nanocage is considered as the least reactive species, it has the lowest value of  $N$ ,  $S$  and the highest  $\chi$ ,  $\eta$  and absolute value of  $\mu$ . Then, it is the most chemically stable species. Due to significant CT from maximum number of titanium and boron heteroatoms to neighboring carbons of C<sub>10</sub>Ti<sub>5</sub>B<sub>5</sub>, it as the most reactive species, shows the lowest  $\mu$  and the highest  $S$  between the scrutinized heterofullerenes.

### 3.4. NBO analysis

Here, NBO analysis is completed on kinetically stable C<sub>18</sub>Ti<sub>1</sub>B<sub>1</sub> nanocage, at the B3PW91/6-311++G\*\* level (Tables 7 and 8).

**Table 7.** The resulted NBO analysis of B3PW91/6-311++G\*\* calculations for the C<sub>18</sub>Ti<sub>1</sub>B<sub>1</sub> heterofullerene.

Bond (A-B)	Occup.	ED <sub>A</sub> (%)	ED <sub>B</sub> (%)	NBO
$\pi_{C=C}$	0.987	50.01	49.99	0.707* $sp^{2.01}$ +0.707* $sp^{2.05}$
$\sigma_{C-B}$	0.976	69.65	30.35	0.834* $sp^{1.52}$ +0.551* $sp^{1.99}$
$\sigma_{Ti-B}$	0.980	46.78	53.22	0.684* $sp^{0.33}d^{4.2}$ +0.729* $sp^{1.78}$
LP(1) <sub>Ti</sub>	0.914			$sp^{0.01}d^{99.99}$

**Table 8.** The hybridization analysis of B3PW91/6-311++G\*\* calculations for the C<sub>18</sub>Ti<sub>1</sub>B<sub>1</sub> heterofullerene.

Bond (A-B)	s (%)	p (%)	d (%)
$\pi_{C=C}$	33.23, 32.77	66.67, 67.15	
$\sigma_{C-B}$	39.58, 33.46	60.34, 66.50	
$\sigma_{Ti-B}$	18.21, 35.98	6.03, 64.00	75.75, 0.02
LP(1) <sub>Ti</sub>	0.00	0.01	99.99

The first two columns give orbital and occupancy between 0.987 electrons for C=C bonding orbitals with  $sp^2$  hybrid vs. 0.914 electrons for LP(1)<sub>Ti</sub> with  $sp^{0.01}d^{99.99}$  hybrid. The latter hybrid on lone pair of Ti has 99.99% d-character. While, occupancy of  $\sigma_{C-B}$  and  $\sigma_{Ti-B}$  is 0.976 and 0.980 electrons, with 0.834\* $sp^{1.52}$ +0.551\* $sp^{1.99}$  and 0.684\* $sp^{0.33}d^{4.2}$ +0.729\* $sp^{1.78}$  hybrids, respectively.

The bonding orbital of  $\sigma_{C-B}$  shows 69.65% carbon character in  $sp^{1.5}$  hybrid and 30.35% boron character in  $sp^{1.99}$  hybrid. Also,  $\sigma_{Ti-B}$  bonding orbital exhibits 46.78% titanium character in  $sp^{0.33}d^{4.2}$  hybrid and 53.22% boron character in  $sp^{1.78}$  hybrid.

Clearly, titanium heteroatom shows lower percentage of NBO and polarization coefficient than the bonding bonds of  $\pi_{C=C}$  and  $\sigma_{C-B}$  because it has lower electronegativity than carbon and boron atoms.

Here, 0.834, 0.551, 0.684 and 0.729 are polarization coefficients.

The magnitude of these coefficients reveals the importance of two hybrids in one bond. The order of polarization coefficients is arranged by carbon more than boron and boron more than titanium. In fact, the intermolecular interaction is formed by the suitable overlap between bonding orbital of  $\pi_{C=C}$  with anti-bonding orbital of LP\*<sub>Ti</sub>, LP\*<sub>B</sub> and  $\sigma^*_{Ti-B}$ , which results



a CT from donating orbital (with a reducing of its occupancy) to accepting orbital (with a rising of its occupancy).

The intermolecular interaction is happened from different hybrids on their atoms with different percentage of s, p and d orbitals. The most important of donor and acceptor is related to  $\pi_{C=C} \rightarrow LP^*(1)_B$ ,  $\pi_{C=C} \rightarrow LP^*(\gamma)_B$ ,  $\pi_{C=C} \rightarrow LP^*(1)_{Ti}$ ,  $\pi_{C=C} \rightarrow LP^*(\gamma)_{Ti}$  and  $\pi_{C=C} \rightarrow \sigma^*_{Ti-B}$  interactions and their corresponding energies is 28.32, 12.53, 4.54, 2.10 and 1.15 kcal/mol, respectively (Table 9).

**Table 9.** The calculated second order perturbation energies  $E^{(2)}$  corresponding to the most important charge transfer interactions (donor-acceptor) of  $C_{18}Ti_1B_1$ , at B3PW91/6-311++G\*\*.

Donor NBO (i)	Acceptor NBO (j)	$E^{(2)}$ (kcal/mol)	$E(j) - E(i)$ (a.u.)	F (i,j) (a.u.)
$\pi_{C=C}$	$LP^*(1)_B$	28.32	0.39	0.42
$\pi_{C=C}$	$LP^*(2)_B$	12.53	0.28	0.078
$\pi_{C=C}$	$LP^*(1)_{Ti}$	4.54	0.21	0.042
$\pi_{C=C}$	$LP^*(2)_{Ti}$	2.10	0.19	0.025
$\pi_{C=C}$	$\sigma^*_{Ti-B}$	1.15	0.34	0.029

#### 4. Conclusion

In this computational research, based on full optimized geometries, none of  $C_{20}$  and its  $C_{20-2n}Ti_nB_n$  derivatives collapses to deformed open-cage as segregated nanostructure. Isolating Ti—B bonds *via* C=C bond or one C atom is a proper approach for achieving substituted heterofullerene such  $C_{10}Ti_5B_5$  since it avoids from five weak homo bonds of Ti—Ti and B—B. The vibrational frequency analysis is applied to assure that all optimized nanocages show no imaginary frequency as true minima. The range of calculated *B.E.* from 12.25 to 38.03 eV/atom reveals the  $C_{20-2n}Ti_nB_n$  nanocages are as more thermodynamically stable heterofullerene and better candidate for hydrogen storage than the Ti-decorated  $B_{38}$  nanofullerene with *B.E.* of 5.67 eV/atom. Here, the absolute heat of atomization of heterofullerenes decreases as number of substituting Ti—B unit increases. Also, DFT calculations based on B3LYP/AUG-cc-pVTZ, M06-2X/6-311++G\*\* and B3PW91/6-311+G\* implies that the band gap as well as conductivity of  $C_{20-2n}Ti_nB_n$  structures is variously depending on number and topology of Ti and B heteroatoms. The obtained NICS (0) data indicates that substituting  $C_{20}$  to  $C_{18}Ti_1B_1$ ,  $C_{16}Ti_2B_2$  and  $C_{14}Ti_3B_3$  leads to considerably more

aromaticity character than benzene, [2.2]paracyclophane and  $C_{20}$  molecules while  $C_{12}Ti_4B_4$  and  $C_{10}Ti_5B_5$  nanocages exhibit anti-aromaticity character. High NBO charge dispersion on  $C_{20-2n}Ti_nB_n$  surfaces with range of -0.401 on carbon atom of  $C_{10}Ti_5B_5$  to +1.084 on titanium atom of  $C_{18}Ti_1B_1$  compared to  $C_{20}$  with range of -0.054 to +0.054, leads to high *B.E.* of  $H_2$  molecule in its absorption on the  $C_{18}Ti_1B_1$  heterofullerene and causes more exploration on its probable appliance for hydrogen storage. The blue colored MEP images reveal the positive charge on Ti sites, red colored MEP images display the negative charge on C and B sites, demonstrating a suitable CT is occurred from Ti substitutions to the neighboring C and B atoms. The  $\langle \alpha \rangle$  value increases from 187.27 a.u. for  $C_{20}$  to 390.42 a.u. for  $C_{12}Ti_4B_4$  species. Obviously, this species exhibits high activity to polar-polar interaction. Also, the global reactivity of  $C_{20}$  and its  $C_{20-2n}Ti_nB_n$  derivatives is increased as the number of substituted Ti and B dopants (n) is increased. Hence,  $C_{18}Ti_1B_1$  as the least chemical reactive species, shows the lowest *N*, *S* and the highest  $\chi$ ,  $\eta$  and absolute value of  $\mu$ . The NBO analysis on  $C_{18}Ti_1B_1$  indicates large CT is take placed from its  $\pi_{C=C}$  bonds to their neighboring  $LP^*_B$ ,  $LP^*_{Ti}$  and  $\sigma^*_{Ti-B}$  anti-bonds.

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